ATOMIC STRUCTURE

- 1. The Nuclear Atom
- 2. Electron Orbits
- 3. Atomic Spectra
- 4. The Bohr Atom
- 5. Energy Level and Spectra
- 6. Correspondence Principle
- 7. Nuclear Motion
- 8. Atomic Excitation
- 9. The Laser

The various permitted orbits involve different electron energies.

$$E = -\frac{e^2}{8\pi\varepsilon_o r}$$

$$E_n = -\frac{e^2}{8\pi\varepsilon_o r_n}$$

$$r_n = \frac{n^2 h^2 \varepsilon_o}{\pi m e^2}$$

Energies.
$$E = -\frac{e^2}{8\pi\varepsilon_o r} \quad E_n = -\frac{e^2}{8\pi\varepsilon_o r_n} \quad r_n = \frac{n^2 h^2 \varepsilon_o}{\pi m e^2}$$

$$E_n = -\frac{me^4}{8\varepsilon_o^2 h^2} \left(\frac{1}{n^2}\right) = \frac{E_I}{n^2}, n = 1, 2, 3, ...$$

$$E_I = -2.18 \times 10^{-18} J = -13.6 \text{ eV}$$

$$E_I = -2.18 \times 10^{-18} J = -13.6 \, eV$$

This equation specifies the energy levels of the H atom.

- →They are negative.
- →They are discrete.
- →Other energies are not allowed.

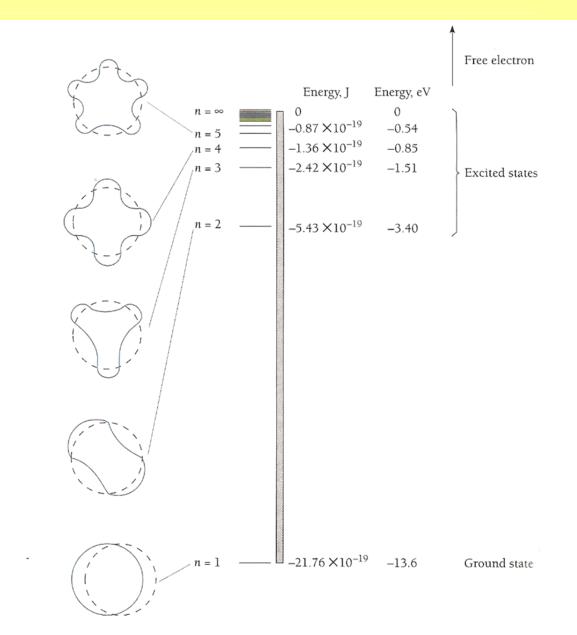
E₁ is called the **ground state**.

 E_2 , E_3 , E_4 ... are called excited states.

 E_{∞} = o \leftarrow electron is not bound to the nucleus.

 E_{T} is the ionization energy.

 $E = +ve \leftarrow$ free electron & has no quantum condition to fulfill.



What is the origin of line spectra and do they follow from Bohr's model?

If the quantum number of the initial (higher-energy) state is n_i and the quantum number of the final (lower-energy) state is n_f , then: initial energy - final energy = emitted photon energy

$$E_i - E_f = h v$$

$$E_{i} - E_{f} = E_{I} \left(\frac{1}{n_{i}^{2}} - \frac{1}{n_{f}^{2}} \right) = -E_{I} \left(\frac{1}{n_{f}^{2}} - \frac{1}{n_{i}^{2}} \right)$$

$$v = \frac{E_i - E_f}{h} = -\frac{E_I}{h} \left(\frac{1}{n_f^2} - \frac{1}{n_i^2} \right)$$

$$\frac{1}{\lambda} = -\frac{E_I}{ch} \left(\frac{1}{n_f^2} - \frac{1}{n_i^2} \right)$$
The Hydrogen spectrum

$$\frac{1}{\lambda} = -\frac{E_I}{ch} \left(\frac{1}{n_f^2} - \frac{1}{n_i^2} \right)$$

$$\frac{1}{\lambda} = -\frac{E_I}{ch} \left(\frac{1}{n_f^2} - \frac{1}{n_i^2} \right)$$

Series		wavelength	
Lyman	n_f =1	$\frac{1}{\lambda} = -\frac{E_I}{ch} \left(\frac{1}{1^2} - \frac{1}{n_i^2} \right)$	n=2,3,4,
Balmer	n_f =2	$\frac{1}{\lambda} = -\frac{E_I}{ch} \left(\frac{1}{2^2} - \frac{1}{n_i^2} \right)$	<i>n</i> =3,4,5,
Paschen	n_f =3	$\frac{1}{\lambda} = -\frac{E_I}{ch} \left(\frac{1}{3^2} - \frac{1}{n_i^2} \right)$	<i>n</i> =4,5,6,
Brackett	n_f =4	$\frac{1}{\lambda} = -\frac{E_I}{ch} \left(\frac{1}{4^2} - \frac{1}{n_i^2} \right)$	<i>n</i> =5,6,7,
Pfund	n_f =5	$\frac{1}{\lambda} = -\frac{E_I}{ch} \left(\frac{1}{5^2} - \frac{1}{n_i^2} \right)$	n=6,7,8,

Can Bohr's model predict the Rydberg constant?

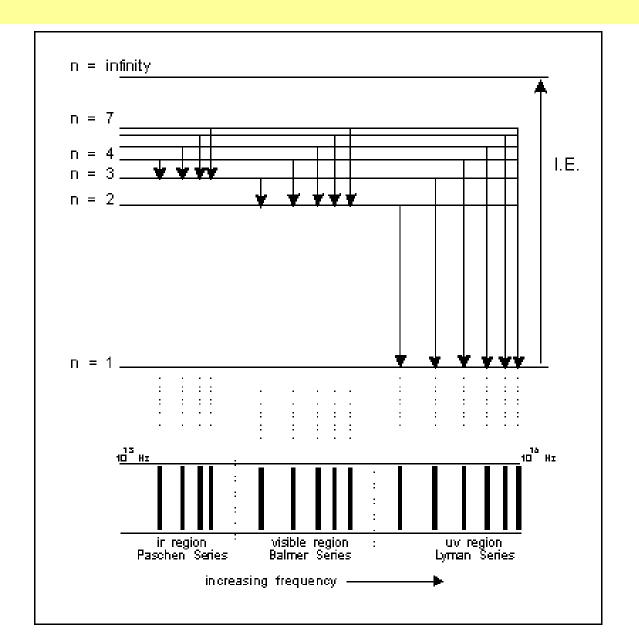
$$\frac{1}{\lambda} = R\left(\frac{1}{2^2} - \frac{1}{n^2}\right) \quad , \quad n = 3,4,5,...$$

$$\frac{1}{\lambda} = -\frac{E_I}{ch}\left(\frac{1}{n_f^2} - \frac{1}{n_i^2}\right)$$

$$R = 1.097 \times 10^7 \ m^{-1}$$

$$\frac{1}{\lambda} = -\frac{E_I}{ch} \left(\frac{1}{n_f^2} - \frac{1}{n_i^2} \right)$$

$$-\frac{E_I}{ch} = \frac{me^4}{8\varepsilon_o^2 ch^3} = 1.097 \times 10^7 m^{-1}$$



Remember...

A photon is emitted when an electron jumps from one energy level to a lower level.

Quantum physics in the microworld must give the same results as classical physics in the macroworld.

How this is applied to the Bohr's model of the H model?

- According to EM theory, an electron moving in a circular orbit radiates EM waves.
- The frequencies of these radiation are equal to its frequency of revolution and to its harmonics. • The electron speed in a H atom is: $v = \frac{e}{\sqrt{4\pi\varepsilon_o mr}} \qquad r_n = \frac{n^2 h^2 \varepsilon_o}{\pi m e^2}$
- The frequency of revolution f of the electron:

$$f = \frac{\text{electron speed}}{\text{orbit circumference}} = \frac{\upsilon}{2\pi r} = \frac{e}{2\pi \sqrt{4\pi\varepsilon_o mr^3}}$$

$$f = \frac{me^4}{8\varepsilon_o^2 h^3} \left(\frac{2}{n^3}\right) = -\frac{E_I}{h} \left(\frac{2}{n^3}\right)$$

Under what circumstances should the Bohr atom behave classically?

• When the electron orbit is so large we might be able to measure it directly. We have:

$$v = -\frac{E_I}{h} \left(\frac{1}{n_f^2} - \frac{1}{n_i^2} \right)$$

Let $n_i \rightarrow n$ & $n_f \rightarrow n - p$ (where p = 1,2,3,...)

$$v = -\frac{E_I}{h} \left[\frac{1}{(n-p)^2} - \frac{1}{n^2} \right] = -\frac{E_I}{h} \left[\frac{2np - p^2}{n^2(n-p)^2} \right]$$

When both $n_i \& n_f$ are very large then n >> p

$$2np - p^{2} \approx 2np$$

$$(n-p)^{2} \approx n^{2}$$

$$v = -\frac{E_{I}}{h} \left(\frac{2p}{n^{3}}\right)$$

$$f = -\frac{E_{I}}{h} \left(\frac{2}{n^{3}}\right)$$

$$v = -\frac{E_I}{h} \left(\frac{2p}{n^3} \right)$$

$$f = -\frac{E_I}{h} \left(\frac{2}{n^3} \right)$$

Under what circumstances should the Bohr atom behave classically?

- The requirement that quantum physics give the same results as classical physics in the limit of large quantum number was called by Bohr the correspondence principle.
- Bohr used the correspondence principle in reverse to look for the condition for orbit stability.
- •A stable orbit must have orbital angular momenta of:

$$m\upsilon r = \frac{nh}{2\pi}$$
 , $n = 1, 2, 3, ...$ $n\lambda = 2\pi r$

Remember...

The greater the quantum number, the closer quantum physics approaches classical physics.